

## CENTRIFUGAL INSTABILITIES IN LID-DRIVEN CAVITY FLOWS

S. Albensoeder<sup>\*1</sup> and H. C. Kuhlmann<sup>1</sup>

<sup>1</sup>ZARM - University of Bremen; Am Fallturm, 28359 Bremen, Germany

### Abstract

The flow of one-sided lid-driven cavities is a fundamental problem in fluid mechanics and has been treated numerically and experimentally by a large number of investigators. In the present paper we numerically calculate, for the first time, the linear stability boundaries of the two-dimensional cavity flow with respect to three-dimensional perturbations for several aspect ratios. Four critical instability branches arise, depending on the aspect ratio. By an energy analysis we claim that all instability branches are of centrifugal type.

### Introduction

The first numerical investigation of the linear stability of the one-sided lid-driven cavity (fig. 1) is due to Ramanan and Homsy [6]. They predicted a critical mode with a wave number  $k \approx 2$  for a cavity with unit aspect ratio. This result was questioned by Kuhlmann *et al.* [4] who provided evidence, using increasing grid sizes, that the neutral stability boundary of this particular mode does not converge to the result of [6]. The first reliable result for a neutral Reynolds number has been given by Ding and Kawahara [2, 3]. They found a short wavelength mode of  $k = 7.4$  and speculated this mode to be critical. More recently, Albensoeder *et al.* [1] also considered modes with higher wave numbers and found that the critical mode has a the wave number  $k = 15.43$ .

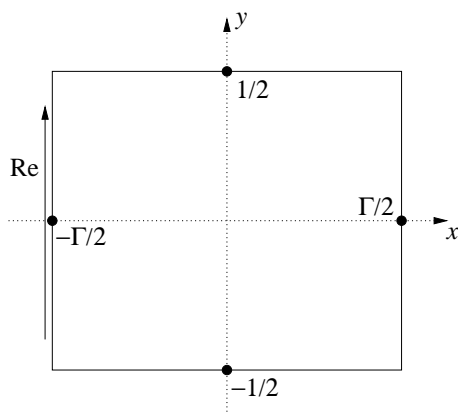


Figure 1: Geometry of the lid-driven cavity.

### Methods of investigation

The basic two-dimensional flow  $(\vec{u}_0, p_0)$  is calculated by solving the steady incompressible Navier-Stokes equations for a viscous fluid with kinematic viscosity  $\nu$ . A finite-volume technique in primitive variables on a staggered grid with refinement near the boundaries is used. The resulting nonlinear equations are solved by Newton-Raphson iteration. To calculate the three-dimensional linear stability we consider normal modes  $(\vec{u}, p) = (\vec{u}_0, p_0) + [\vec{u}(x, y), \hat{p}(x, y)]e^{\sigma t + i(kz - \omega t)}$ . The resulting generalized eigenvalue problem is solved by inverse iteration. While the typical grid resolution is  $141 \times 141$ , convergence test were made by comparison with  $71 \times 71$  grid points. Lengths, velocities, pressure, and time are scaled by  $h, \nu/h, \rho\nu^2/h^2$ , and  $h^2/\nu$ , respectively, where  $h$  is the height of the square cavity. The non-dimensional parameters of the problem are the Reynolds number  $Re = Vh/\nu$  and the aspect ratio  $\Gamma = d/h$  with  $V$  being the lid velocity and  $d$  the width of the cavity.

### Results

Figure 2 shows the calculated neutral Reynolds numbers  $Re_n$ , wave numbers  $k$ , and frequencies  $\omega$  as functions of the aspect ratio  $\Gamma$ . It turns out that four different modes may become critical, depending on the aspect ratio. Two modes in the intervals  $\Gamma \in [0, 0.888]$  and  $\Gamma \in [1.163, 1.207]$ , respectively, are oscillatory, while the other neutral modes are steady. The critical mode for  $\Gamma \rightarrow 0$  scales asymptotically like  $Re_c^* = \Gamma Re_c(\Gamma \rightarrow 0) \approx 275$ ,  $k_c^* = \Gamma k_c(\Gamma \rightarrow 0) \approx 5.1$  and  $\omega_c^* = \Gamma^2 \omega_c(\Gamma \rightarrow 0) \approx 124$ , where  $Re_c^*, k_c^*$ ,

and  $\omega_c^*$  represent quantities scaled with the width  $d$  instead of the height  $h$ .

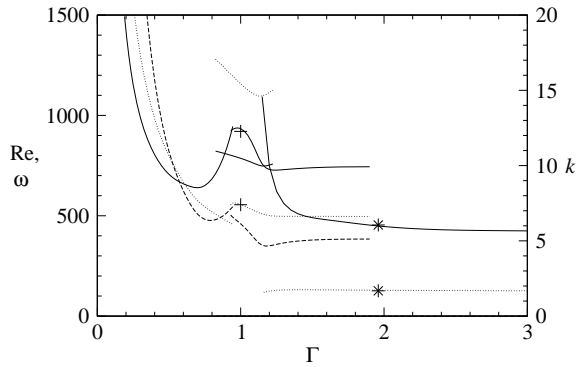


Figure 2: Neutral Reynolds numbers  $Re_n$  (solid lines) as functions of the aspect ratio  $\Gamma$ . The dotted and dashed lines denotes the wave number  $k$  and the oscillation frequency  $\omega$  of the neutral mode, respectively. Previous results of [3] and [4, 5] are indicated by (+) and (\*), respectively.

For  $\Gamma \in [0.888, 1.163]$  and  $\Gamma \in [1.207, \infty]$  the critical modes are stationary. The asymptotic critical Reynolds number and wave number for infinitely wide cavities are  $Re_c(\Gamma \rightarrow \infty) \approx 420$  and  $k_c(\Gamma \rightarrow \infty) \approx 1.685$ .

For a detailed analysis of the instability process an *a posteriori* energy analysis [1] has been carried out. It turned out that all instabilities are caused by centrifugal effects. This conclusion is supported by application of the criterion of Sipp and Jacquin [7] for centrifugal instabilities.

## Conclusion

The linear stability problem of the two-dimensional flow in one-sided lid-driven cavities has been solved for various aspect ratios. All critical modes are of centrifugal type. The asymptotic scaling for  $\Gamma \rightarrow 0$  and  $\Gamma \rightarrow \infty$  has been established..

## Acknowledgement

This work has been supported by DFG under grant numbers Ku896/5-2 and Ku896/8-1.

## References

- [1] Albensoeder, S., Kuhlmann, H. C. and Rath, H. J., 2001, "Three-Dimensional Centrifugal-Flow Instabilities in the Lid-Driven Cavity Problem", *Phys. Fluids*, 13:121–135.
- [2] Ding, Y. and Kawahara, M., 1998, "Linear stability of incompressible fluid flow in a cavity

using finite element method", *Int. J. Numer. Meth. Fluids*, 27:139–157.

- [3] Ding, Y. and Kawahara, M., 1999, "Three-dimensional linear stability analysis of incompressible viscous flows using the finite element method", *Int. J. Numer. Meth. Fluids*, 31:451–479.
- [4] Kuhlmann, H. C., Wanschura, M. and Rath, H. J., 1997, "Flow in two-sided lid-driven cavities: Non-uniqueness, instabilities, and cellular structures", *J. Fluid Mech.*, 336:267–299.
- [5] Kuhlmann, H. C., Wanschura, M. and Rath, H. J., 1998, "Elliptic instability in two-sided lid-driven cavity flow", *Eur. J. Mech., B/Fluids*, 17:561–569.
- [6] Ramanan, N. and Homsy, G. M., 1994, "Linear stability of lid-driven cavity flow", *Phys. Fluids*, 8:2690–2701.
- [7] Sipp, D. and Jacquin, L., 2000, "Three-dimensional centrifugal-type instabilities of two-dimensional flows in rotating systems", *Phys. Fluids*, 12:1740–1748.